

# Generalized Cat States

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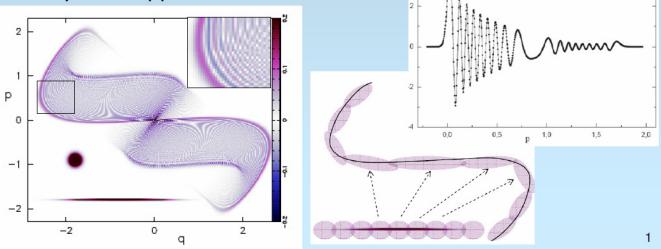
## Motivation

$$W^\Psi(x) \propto \langle \Psi | \hat{R}_x | \Psi \rangle$$

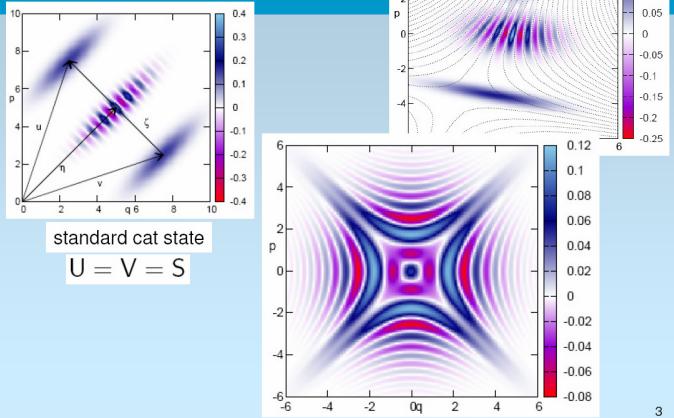
Pure superposition:  $|\Psi\rangle = \sum_{i=1}^M C_i |\psi_i\rangle$ ,

$$W^\Psi(x) := \frac{1}{(\pi\hbar)^n} \left[ \sum_{i=1}^M |C_i|^2 \langle \psi_i | \hat{R}_x | \psi_i \rangle \right] + \frac{2}{(\pi\hbar)^n} \operatorname{Re} \left[ \sum_{j>i=1}^M C_i C_j^* \langle \psi_j | \hat{R}_x | \psi_i \rangle \right]$$

Nearby orbit approximation:



## 2D Gallery



## Decoherence of Generalized Cats

$$\frac{\partial \hat{\rho}}{\partial t} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] + \frac{1}{\hbar} \sum_k \left( \hat{L}_k \hat{\rho} \hat{L}_k^\dagger - \frac{1}{2} \hat{L}_k^\dagger \hat{L}_k \hat{\rho} - \frac{1}{2} \hat{\rho} \hat{L}_k^\dagger \hat{L}_k \right),$$

linear Lindblads:  $\hat{L}_j(\hat{x}) := l'_j \cdot \hat{x} + i l''_j \cdot \hat{x}$  + quadratic H

$$\frac{\partial W}{\partial t}(x, t) = \mathcal{L}_W^H [W(x, t)] + \mathcal{L}_W^b [W(x, t)] + \mathcal{L}_W^\# [W(x, t)]$$

$$\mathcal{L}_W^H [W(x, t)] = \{H(x), W(x, t)\} \longrightarrow \text{Hamiltonian}$$

$$\mathcal{L}_W^b [W(x, t)] = \frac{\hbar}{2} \sum_{j=1}^M \left[ J'_j \cdot \frac{\partial^2 W}{\partial x \partial x} J'_j + J''_j \cdot \frac{\partial^2 W}{\partial x \partial x} J''_j \right] \longrightarrow \text{diffusive}$$

$$\mathcal{L}_W^\# [W(x, t)] = \alpha \left[ x \cdot \frac{\partial W}{\partial x} + 2nW(x, t) \right] \xrightarrow{\alpha := \sum_{j=1}^M l'_j \wedge l''_j} \text{contractive}$$

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## Wigner Function

$$|S; \xi\rangle := \hat{T}_\xi \hat{M}_S |0\rangle$$

$$2\text{-Superposition: } |\psi\rangle = \frac{1}{\sqrt{1+2\operatorname{Re}[a^*b\langle U; u|V; v\rangle]}} (a|U; u\rangle + b|V; v\rangle)$$

Wigner:

$$W_\psi(x) = \frac{1}{1+2\operatorname{Re}[a^*b\langle U; u|V; v\rangle]} \left\{ |a|^2 \mathcal{G}_{2n}^R[x; (UU^\top)^{-1}; u] + |b|^2 \mathcal{G}_{2n}^R[x; (VV^\top)^{-1}; v] \right\}$$

$$\xrightarrow{\text{interference}} + 2 \operatorname{Re}[a^*b\langle U; u|\hat{R}_x|V; v\rangle]$$

$$\langle 0 | \hat{M}_U^\dagger \hat{T}_u^\dagger \hat{R}_x \hat{T}_v \hat{M}_V | 0 \rangle = \frac{2^n i^{\nu - \frac{1}{2} \operatorname{sng} C_0^{-1} V}}{\sqrt{\det[(U+V) + i(U-V)J]}} e^{\frac{i}{\hbar} x \wedge \zeta + \frac{i}{2\hbar} \zeta \wedge \bar{\eta}} \mathcal{G}_{2n}^C[x; G; \bar{\eta}]$$

$$G = 2(UU^\top + VV^\top)^{-1} - i(UU^\top + VV^\top)^{-1}(UU^\top - VV^\top) J$$

$$\bar{\eta} = \frac{1}{2}(u+v) \quad e \quad \zeta := u-v \quad \text{Complex, Symmetric, Symplectic}$$

## 2N Dimensions: Euler Decomposition & Sylvester Law of Inertia

$$\text{Euler decomposition for } G' = 2(\mathbb{I}_2 + VV^\top)^{-1} - i(\mathbb{I}_2 + VV^\top)^{-1}(\mathbb{I}_2 - VV^\top) J.$$

$$S = OZO', \quad \text{com} \begin{cases} Z = \operatorname{Diag}(z_1, \dots, z_n, 1/z_1, \dots, 1/z_n), & z_i \geq 1 \ \forall i \\ O', O \in \operatorname{Sp}(2n, \mathbb{R}) \cap \operatorname{SO}(2n) \end{cases}$$

$$SS^\top = OZO' O'^\top ZO^\top = OZ^2 O^\top$$

$$\text{law of Inertia for } G = 2(UU^\top + VV^\top)^{-1} - i(UU^\top + VV^\top)^{-1}(UU^\top - VV^\top) J$$

$$(n_+, n_-, n_0) [BAB^\top] = (n_+, n_-, n_0) [A]$$

In 2N-D the intereference pattern is also hyperbolic !!

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## Decoherence II

Lie-Trotter

Gaussian case

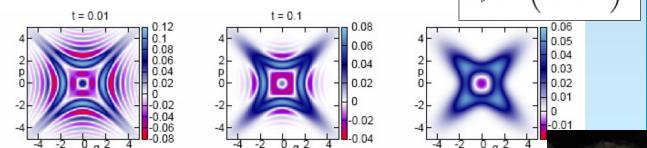
$$W(x, t) = \mathcal{G}_{2n}^R[x; e^{2\gamma t} (\mathcal{S}\Delta_0 S^\top + \Gamma)^{-1} ; e^{-\gamma t} S \eta_0]$$

$S \rightarrow$  Hamiltonian !  
 $\gamma \rightarrow$  contraction !  
 $\Gamma \rightarrow$  diffusion !?

$$\gamma_t := \int_0^t \alpha dt' \quad \Gamma = 2 \int_0^t \sum_{j=1}^M J^\top (l'_j l'_j{}^\top + l''_j l''_j{}^\top) J d\tau$$

$$\mathcal{I}_b(x) = \frac{2^n i^{\nu - \frac{1}{2} \operatorname{sng} C_0^{-1} V}}{\sqrt{\det[(U+V) + i(U-V)J]}} e^{-\frac{1}{4\hbar} \zeta \wedge \Gamma \Omega_b \zeta - \frac{i}{2\hbar} \bar{\eta} \cdot [G\Gamma - I_{2n}] \Omega_b J \zeta} \times \mathcal{G}_{2n}^C[x; E_b; \bar{\eta}_b] \exp(-\frac{i}{\hbar} x \cdot \Omega_b J \zeta),$$

$$E_b := \left( \mathcal{G}^{-1} + \Gamma \right)^{-1}$$



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